

## Topics Summary for Final

### 1. Quantum Spin States

#### a. State Vectors

$$i. |\psi\rangle = c_+ |+\rangle + c_- |-\rangle = |+\rangle \langle +|\psi\rangle + |-\rangle \langle -|\psi\rangle$$

$$ii. \langle \psi| = c_+^* \langle +| + c_-^* \langle -| = \langle \psi|+\rangle \langle +| + \langle \psi|-\rangle \langle -|$$

#### b. Inner Products:

$$i. \langle +|+\rangle = \langle -|-\rangle = 1, \text{ normalized, } \langle +|-\rangle = \langle -|+\rangle = 0. \text{ orthogonal}$$

$$ii. \langle \psi|\psi\rangle = (c_+^* \langle +| + c_-^* \langle -|)(c_+ |+\rangle + c_- |-\rangle) = |c_+|^2 + |c_-|^2$$

$$iii. |\langle \phi|\psi\rangle|^2 = \text{Probability that a particle in state } |\psi\rangle \text{ can be found in state } |\phi\rangle.$$

#### c. States in $S_z$ -basis

$$i. |\pm x\rangle = \frac{1}{\sqrt{2}} |+\rangle \pm \frac{1}{\sqrt{2}} |-\rangle, \langle \pm x| = \frac{1}{\sqrt{2}} \langle +| \pm \frac{1}{\sqrt{2}} \langle -|$$

$$|\pm y\rangle = \frac{1}{\sqrt{2}} |+\rangle \pm \frac{i}{\sqrt{2}} |-\rangle, \langle \pm y| = \frac{1}{\sqrt{2}} \langle +| \mp \frac{i}{\sqrt{2}} \langle -|$$

#### ii. General States

$$1. |\psi\rangle = \sum_n c_n |a_n\rangle, \langle \psi| = \sum_n c_n^* \langle a_n|$$

$$2. \langle \psi|\psi\rangle = \sum_{i,j} c_i^* c_j \langle a_i|a_j\rangle = \sum_j |c_j|^2$$

$$iii. \text{Expectation value: } \langle A \rangle = \sum_n P(a_n) a_n = \sum_n |c_n|^2 a_n$$

$$iv. \text{Uncertainty: } \Delta A = \sqrt{\langle (A - \langle A \rangle)^2 \rangle} = \sqrt{\langle A^2 \rangle - \langle A \rangle^2}$$

#### d. Matrix Representations in $S_z$ -basis

$$i. |\psi\rangle = c_+ |+\rangle + c_- |-\rangle \rightarrow \begin{pmatrix} c_+ \\ c_- \end{pmatrix} = \begin{pmatrix} \langle +|\psi\rangle \\ \langle -|\psi\rangle \end{pmatrix}$$

$$ii. \langle \psi| = c_+^* \langle +| + c_-^* \langle -| \rightarrow (c_+^*, c_-^*) = (\langle \psi|+\rangle, \langle \psi|-\rangle)$$

### 2. Operators

#### a. Rotations

$$i. \hat{R}(d\phi \mathbf{k}) = 1 - \frac{i}{\hbar} \hat{J}_z d\phi, \hat{J}_z |\pm z\rangle = \pm \frac{\hbar}{2} |\pm z\rangle$$

$$ii. \hat{R}(\phi \mathbf{k}) = e^{-i\hat{J}_z \phi / \hbar}$$

$$iii. |+\rangle = \hat{R}\left(\frac{\pi}{2} \mathbf{j}\right) |+\rangle, \langle +x| = \langle +z| \hat{R}^\dagger\left(\frac{\pi}{2} \mathbf{j}\right)$$

#### b. Projections (Spin 1/2)

$$i. \hat{P}_+ = |+\rangle \langle +|, \hat{P}_- = |-\rangle \langle -|, |+\rangle \langle +| + |-\rangle \langle -| = I$$

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$$\text{ii. } \hat{P}_{\pm}^2 = \hat{P}_{\pm}, \hat{P}_{+}\hat{P}_{-} = \hat{P}_{-}\hat{P}_{+} = 0$$

c. Matrix Representations of Operators,  $A_{ij} = \langle i | \hat{A} | j \rangle$

d. Commutators

$$\text{i. } [\hat{J}_x, \hat{J}_y] = i\hbar\hat{J}_z, [\hat{J}_y, \hat{J}_z] = i\hbar\hat{J}_x, [\hat{J}_z, \hat{J}_x] = i\hbar\hat{J}_y$$

ii. Commuting operators have simultaneous eigenstates.

3. Matrix/Operator Types

a. Adjoint  $A^\dagger = (A^t)^*$ , Unitary  $A^\dagger A = I$ , Hermitian  $A^\dagger = A$

4. Operator Relations

$$\text{a. } \hat{J}^2 = \hat{J}_x^2 + \hat{J}_y^2 + \hat{J}_z^2$$

$$\text{b. } \hat{J}_{\pm} = \hat{J}_x \pm i\hat{J}_y$$

$$\text{c. } \hat{J}_-\hat{J}_+ = \hat{J}^2 - \hat{J}_z^2 - \hbar\hat{J}_z, \hat{J}_+\hat{J}_- = \hat{J}^2 - \hat{J}_z^2 + \hbar\hat{J}_z$$

5. Eigenvalues and Eigenstates

$$\text{a. } \hat{J}^2 |j, m\rangle = j(j+1)\hbar^2 |j, m\rangle, \hat{J}_z |j, m\rangle = m\hbar |j, m\rangle$$

$$\text{b. } m = -j, -j+1, \dots, j-1, j$$

6. Raising and Lowering Operators

$$\text{a. } \hat{J}_{\pm} |j, m\rangle = \sqrt{j(j+1) - m(m \pm 1)} \hbar |j, m \pm 1\rangle$$

7. Matrix Representations (Note similarity between  $J$ 's and  $S$ 's)

$$\text{a. } \hat{A} |\psi\rangle = |\phi\rangle \rightarrow \begin{pmatrix} A_{11} & A_{12} \\ A_{21} & A_{22} \end{pmatrix} \begin{pmatrix} \langle 1 | \psi \rangle \\ \langle 2 | \psi \rangle \end{pmatrix} = \begin{pmatrix} \langle 1 | \phi \rangle \\ \langle 2 | \phi \rangle \end{pmatrix}$$

$$\text{b. } \hat{P}_+ = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}, \hat{P}_- = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\text{c. } \hat{R}(\phi \mathbf{k}) = \begin{pmatrix} e^{-i\phi/2} & 0 \\ 0 & e^{i\phi/2} \end{pmatrix}$$

8. Change of Basis

$$\text{a. } |\psi'\rangle \xrightarrow{S_z \text{ Basis}} \begin{pmatrix} \langle +z | \psi' \rangle \\ \langle -z | \psi' \rangle \end{pmatrix} = \begin{pmatrix} \langle +z | \hat{R}^\dagger | \psi \rangle \\ \langle -z | \hat{R}^\dagger | \psi \rangle \end{pmatrix} = \begin{pmatrix} \langle +x | \psi \rangle \\ \langle -x | \psi \rangle \end{pmatrix} \xleftarrow{S_x \text{ Basis}} |\psi\rangle$$

$$\text{b. } \hat{A} \xrightarrow{S_z \text{ basis}} A_z = \begin{pmatrix} \langle +z | \hat{A} | +z \rangle & \langle +z | \hat{A} | -z \rangle \\ \langle -z | \hat{A} | +z \rangle & \langle -z | \hat{A} | -z \rangle \end{pmatrix}$$

$$\text{c. } \hat{A} \rightarrow A_x = S^\dagger A_z S, \text{ where } S = \begin{pmatrix} \langle +z | +x \rangle & \langle +z | -x \rangle \\ \langle -z | +x \rangle & \langle -z | -x \rangle \end{pmatrix}$$

d. Spin  $\frac{1}{2}$  -  $\hat{S} \rightarrow \frac{\hbar}{2} \boldsymbol{\sigma}$  in terms of Pauli spin matrices.

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$$\hat{S}_x = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \hat{S}_y = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \hat{S}_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

e. Matrix Elements:  $(A_{mm'}) = (\langle j, m' | \hat{A} | j, m \rangle)$  where  $m$  is the row,  $m'$  is the column and one proceeds in the order  $m = j, j-1, \dots, -j+1, -j$ .

f. 3D Rotation matrices

$$g. \hat{S}(\phi\mathbf{i}) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\phi & -\sin\phi \\ 0 & \sin\phi & \cos\phi \end{pmatrix}, \hat{S}(\phi\mathbf{j}) = \begin{pmatrix} \cos\phi & 0 & \sin\phi \\ 0 & 1 & 0 \\ -\sin\phi & 0 & \cos\phi \end{pmatrix}, \hat{S}(\phi\mathbf{k}) = \begin{pmatrix} \cos\phi & -\sin\phi & 0 \\ \sin\phi & \cos\phi & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

h. Spin 1

$$S_x = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}, S_y = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -i & 0 \\ i & 0 & -i \\ 0 & i & 0 \end{pmatrix}, S_z = \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix},$$

$$S_+ = \sqrt{2}\hbar \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{pmatrix}, S_- = \sqrt{2}\hbar \begin{pmatrix} 0 & 0 & 0 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}, S^2 = 2\hbar^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$i. \text{ Spin-1 Representation: } A \rightarrow \begin{pmatrix} \langle 1,1|A|1,1\rangle & \langle 1,1|A|1,0\rangle & \langle 1,1|A|1,-1\rangle \\ \langle 1,0|A|1,1\rangle & \langle 1,0|A|1,0\rangle & \langle 1,0|A|1,-1\rangle \\ \langle 1,-1|A|1,1\rangle & \langle 1,-1|A|1,0\rangle & \langle 1,-1|A|1,-1\rangle \end{pmatrix}$$

9. Uncertainty  $\Delta A \Delta B \geq \frac{|\langle C \rangle|}{2}$  for  $[A, B] = iC$ , and  $A, B, C$  Hermitian

10. Send particles with spin  $s$  through Stern-Gerlach device oriented in different positions.

Determine the number of beam channels and the probability a given particle will be found in a given channel.

11. Time Evolution:

$$a. |\psi(t)\rangle = \hat{U}(t) |\psi(0)\rangle = e^{-i\hat{H}t/\hbar} |\psi(0)\rangle$$

$$b. i\hbar \frac{d}{dt} |\psi(t)\rangle = \hat{H}(t) |\psi(t)\rangle, \text{ Schrödinger Equation}$$

$$c. \frac{d}{dt} \langle A \rangle = \langle \psi(t) | \left( \frac{i}{\hbar} [\hat{H}, \hat{A}] + \frac{\partial \hat{A}}{\partial t} \right) | \psi(t) \rangle$$

$$12. \text{ Hamiltonians: } \hat{H} = -\hat{\mu} \cdot \mathbf{B} = \omega_0 \hat{S}_z, \hat{H} = \frac{2A}{\hbar^2} \hat{\mathbf{S}}_1 \cdot \hat{\mathbf{S}}_2$$

$$13. \text{ Precession of Spin: } \langle S_z \rangle = 0, \langle S_x \rangle = \frac{\hbar}{2} \cos \omega_0 t, \langle S_y \rangle = \frac{\hbar}{2} \sin \omega_0 t, \omega_0 = \frac{ge}{2mc} B_0$$

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### 14. Addition of Angular Momenta, Eigenstates of Total Angular Momenta – two particles

- $R(d\theta\mathbf{n}) = 1 - \frac{i}{\hbar} \hat{\mathbf{S}} \cdot \mathbf{n} d\theta = \left(1 - \frac{i}{\hbar} \hat{\mathbf{S}}_1 \cdot \mathbf{n} d\theta\right) \otimes \left(1 - \frac{i}{\hbar} \hat{\mathbf{S}}_2 \cdot \mathbf{n} d\theta\right)$
- $\hat{\mathbf{S}} = \hat{\mathbf{S}}_1 \otimes 1 + 1 \otimes \hat{\mathbf{S}}_2$
- $\hat{\mathbf{S}}^2 = (\hat{\mathbf{S}}_1 + \hat{\mathbf{S}}_2)^2 = \hat{\mathbf{S}}_1^2 + \hat{\mathbf{S}}_2^2 + 2\hat{\mathbf{S}}_1 \cdot \hat{\mathbf{S}}_2, \hat{S}_z = \hat{S}_{1z} + \hat{S}_{2z}$

### 15. Continuous States

- $|\psi\rangle = \int da |a\rangle \langle a|\psi\rangle, \langle a'|a\rangle = \delta(a'-a)$
- $\hat{x}|x\rangle = x|x\rangle, \hat{p}_x|p\rangle = p|p\rangle$
- $\hat{T}(a) = e^{-i\hat{p}_x a/\hbar}, \hat{T}(a)|x\rangle = |x+a\rangle$
- $\langle x|\hat{p}_x|\psi\rangle = \frac{\hbar}{i} \frac{\partial}{\partial x} \langle x|\psi\rangle$
- $\langle p|\psi\rangle = \int dx \langle p|x\rangle \langle x|\psi\rangle = \frac{1}{\sqrt{2\pi\hbar}} \int e^{-ipx/\hbar} \langle x|\psi\rangle dx$

### 16. Schrödinger's Equation

- $\langle x|\hat{H}|\psi\rangle = \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x)\right] \langle x|\psi\rangle = i\hbar \frac{\partial}{\partial t} \langle x|\psi\rangle$
- 1D Time Independent:  $-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V(x)\psi = E\psi$
- 3D Time Independent:  $-\frac{\hbar^2}{2m} \nabla^2\psi + V(x)\psi = E\psi$

### 17. Energy eigenstates: $\hat{H}|E_n\rangle = E_n|E_n\rangle, |\psi\rangle = \sum_n \langle E_n|\psi\rangle |E_n\rangle$

### 18. One Dimensional Problems

- Ehrenfest's Theorem  $\frac{d\langle x\rangle}{dt} = \frac{\langle p_x\rangle}{m}, \frac{d\langle p_x\rangle}{dt} = \left\langle -\frac{dV}{dx} \right\rangle \Rightarrow m \frac{d^2\langle x\rangle}{dt^2} = \left\langle -\frac{dV}{dx} \right\rangle$
- Free particle – definition and time evolution
- Gaussian wave packet -  $\int e^{-ipx/\hbar - \alpha x^2} dx$  - Computing Gaussian integrals  $\int_{-\infty}^{\infty} x^n e^{-\beta x^2} dx,$   
 $\langle x\rangle, \langle x^2\rangle, \Delta x, \langle p_x\rangle, \langle p_x^2\rangle, \Delta p_x,$  and uncertainty  $\Delta x \Delta p_x$
- Particle in Box, Infinite Square Well,  $0 \leq x \leq a.$ 
  - $\psi_n(x) = \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a}, E_n = \frac{\hbar^2 \pi^2 n^2}{2ma^2}$
- Finite Square Well
- Bound states and scattering states
- Reflection, Transmission probabilities

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### h. Tunneling

### i. Harmonic Oscillator

- i.  $\hat{H} = \left( \hat{a}^\dagger \hat{a} + \frac{1}{2} \right) \hbar \omega, [\hat{a}, \hat{a}^\dagger] = 1$
- ii.  $\hat{a}^\dagger |n\rangle = \sqrt{n+1} |n+1\rangle, \hat{a} |n\rangle = \sqrt{n} |n-1\rangle$
- iii.  $\hat{x} = \sqrt{\frac{\hbar}{2m\omega}} (\hat{a} + \hat{a}^\dagger), \hat{p}_x = -i\sqrt{\frac{m\omega\hbar}{2}} (\hat{a} - \hat{a}^\dagger)$
- iv.  $-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + \frac{1}{2} m\omega^2 x^2 \psi = E\psi$
- v.  $E_n = \left( n + \frac{1}{2} \right) \hbar \omega$
- vi.  $\langle x|n\rangle = \frac{1}{\sqrt{n!}} \langle x|(\hat{a}^\dagger)^n|0\rangle, \langle x|0\rangle = \left( \frac{m\omega}{\pi\hbar} \right)^{\frac{1}{4}} e^{-m\omega x^2/2\hbar}$

## 19. Two-body Problem

- a.  $\hat{H}|E, \ell, m\rangle = E|E, \ell, m\rangle$
- b.  $\hat{L}^2|E, \ell, m\rangle = \ell(\ell+1)\hbar^2|E, \ell, m\rangle$
- c.  $\hat{L}_z|E, \ell, m\rangle = m\hbar|E, \ell, m\rangle$
- d.  $[\hat{L}_x, \hat{L}_y] = i\hbar\hat{L}_z, \text{ etc.}$
- e.  $\hat{L}_z \rightarrow \frac{\hbar}{i} \frac{\partial}{\partial \varphi}, \hat{L}_x \rightarrow \frac{\hbar}{i} \left( -\sin\varphi \frac{\partial}{\partial \theta} - \cot\theta \cos\varphi \frac{\partial}{\partial \varphi} \right), \hat{L}_y \rightarrow \frac{\hbar}{i} \left( \cos\varphi \frac{\partial}{\partial \theta} - \cot\theta \sin\varphi \frac{\partial}{\partial \varphi} \right)$
- f. Raising and lowering operators:  $\hat{L}_\pm \rightarrow \frac{\hbar}{i} e^{\pm i\varphi} \left( \pm i \frac{\partial}{\partial \theta} - \cot\theta \frac{\partial}{\partial \varphi} \right)$  - applied to  $\langle \theta, \varphi | \ell, \ell \rangle = c_\ell e^{i\ell\varphi} \sin^\ell \theta, L_- | \ell, m \rangle = \sqrt{\ell(\ell+1) - m(m-1)} \hbar | \ell, m-1 \rangle$ . Here  $c_\ell = \frac{(-1)^\ell}{2^\ell \ell!} \sqrt{\frac{(2\ell+1)!}{4\pi}}$ .
- g. Spherical Harmonics  $Y_{\ell m}(\theta, \varphi) = P_\ell^m(\cos\theta) e^{\pm im\varphi}$
- h. Hydrogen Atom

- i.  $-\frac{\hbar^2}{2\mu} \nabla^2 \psi + \left[ \frac{\ell(\ell+1)\hbar^2}{2\mu r^2} - \frac{Zke^2}{r} \right] \psi = E\psi$
- ii.  $\psi(r, \theta, \varphi) = R(r) Y_\ell^m(\theta, \varphi)$
- iii.  $E = -\frac{\mu k^2 e^4}{2\hbar^2 n^2} Z^2 = -\frac{\mu c^2 \alpha^2}{2n^2} Z^2 = -13.6 \text{ eV} \frac{Z}{n^2}$

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$$\text{iv. } a_0 = \frac{\hbar}{\mu c \alpha} = 0.529 \text{ \AA}, \alpha = \frac{ke^2}{\hbar c} \approx \frac{1}{137.036}$$

i. Deuteron Model,  $R(r) = ru(r)$

$$\text{i. } -\frac{\hbar^2}{2\mu} u'' - V_0 u = Eu, r < a \Rightarrow \tan k_0 a = -\frac{k_0}{q}$$

j. Infinite Square Well, spherical Bessel functions, and Magic Numbers

20. Time Independent Perturbation Theory

$$\text{a. } \hat{H} = \hat{H}_0 + \lambda \hat{H}_1, \hat{H}_0 |\phi_n^{(0)}\rangle = E_n^{(0)} |\phi_n^{(0)}\rangle$$

$$\text{b. Energy Shift: } E_n^{(1)} = \langle \phi_n^{(0)} | \hat{H}_1 | \phi_n^{(0)} \rangle,$$

$$\text{c. Wave function: } |\phi_n^{(1)}\rangle = \sum_{k \neq n} |\phi_k^{(0)}\rangle \frac{\langle \phi_k^{(0)} | \hat{H}_1 | \phi_n^{(0)} \rangle}{E_n^{(0)} - E_k^{(0)}}$$

$$\text{d. Degenerate States, } (H_1)_{ji} = \langle \phi_{n,j}^{(0)} | \hat{H}_1 | \phi_{n,i}^{(0)} \rangle$$

$$\text{i. } \begin{pmatrix} (H_1)_{11} & (H_1)_{12} \\ (H_1)_{21} & (H_1)_{22} \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = E_n^{(1)} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix}$$

$$\text{e. Stark Effect, } \hat{H}_1 = -\hat{\mu}_e \cdot \mathbf{E} = e\mathbf{r} \cdot \mathbf{E}$$

$$\text{f. Spin-orbit Coupling } \hat{H}_1 = -\hat{\mu} \cdot \mathbf{B} = \frac{Zke^2}{m^2 c^2 r^3} \mathbf{S} \cdot \mathbf{L}$$

g. Darwin Term, Relativistic Correction

h. Fine and hyperfine structure, fine structure constant

$$\text{i. Zeeman Effect, } \hat{H}_1 = -\hat{\mu} \cdot \mathbf{B} = \left( \frac{e}{2mc} \mathbf{L} + \frac{e}{mc} \mathbf{S} \right) \cdot \mathbf{B}$$

j. Perturbations of: particle in a box (1D, 2D, 3D), harmonic oscillators (1D, 2D, 3D), Spin systems (Spin  $\frac{1}{2}$ , Spin 1, etc.)