Fourier-Bessel Series Examples

We consider examples of the Fourier-Bessel function series on [0, b]

$$f(x) = \sum_{i=1}^{\infty} c_i J_n(\alpha_i x), \tag{1}$$

Spring 2025

Revised: April 21, 2025

where

$$c_i = \frac{\int_0^b x J_n(\alpha_i x) f(x) dx}{\|J_n(\alpha_i x)\|^2}, \quad i = 1, 2, \dots$$
 (2)

Here, $J_n(\alpha x)$ is a solution of Bessel's equation in the form

$$\frac{d}{dx}\left(x\frac{dy}{dx}\right) + \left(\alpha^2 x - \frac{n^2}{x}\right)y = 0. \tag{3}$$

We assume that one boundary condition is that the solution is finite at the origin. The other homogeneous boundary conditions at x = b take the form

$$Ay(b) + By'(b) = 0.$$

Noting that from the Chain Rule we have

$$\frac{dJ_n(\alpha x)}{dx} = \alpha J_n'(\alpha x),$$

then this boundary condition can be written as

$$AJ_n(\alpha b) + \alpha BJ'_n(\alpha b) = 0.$$

Solutions of this equation yield the values α_i , for $i = 1, 2, \ldots$ This also impacts the value of $||J_n(\alpha_i x)||^2$ as well as providing the orthogonality condition

$$\int_0^b x J_n(\alpha_i x) J_n(\alpha_j x) dx = ||J_n(\alpha_i x)||^2 \delta_{ij}.$$

One can derive the expansion coefficients for different types of boundary conditions. The main affect is in the value of $||J_n(\alpha_i x)||^2$.

• For $J_n(\alpha b) = 0$, we have that

$$||J_n(\alpha_i x)||^2 = \frac{b^2}{2} [J_{n+1}(\alpha_i b)]^2.$$

• For $J_0'(\alpha b) = 0$, we have that

$$||J_0(\alpha_i x)||^2 = \frac{b^2}{2} [J_0(\alpha_i b)]^2,$$

for i > 1. For i = 1, we have

$$||J_0(\alpha_i x)||^2 = \frac{b^2}{2}.$$

• For $AJ_n(\alpha b) + \alpha BJ'_n(\alpha b) = 0$, we have that

$$||J_n(\alpha_i x)||^2 = \frac{\alpha_i^2 b^2 - n^2 + A^2}{2\alpha_i^2} [J_n(\alpha_i b)]^2.$$

We now look at a couple of examples of Fourier-Bessel series. We note that several identities may be useful.

$$\frac{d}{dx}\left[x^n J_n(x)\right] = x^n J_{n-1}(x),\tag{4}$$

$$\frac{d}{dx} \left[x^{-n} J_n(x) \right] = x^{-n} J_{n+1}(x). \tag{5}$$

Example 1. Consider the function $f(x) = x^2$, $x \in [0,1]$. Find the Fourier-Bessel series expansion in $J_2(\alpha_i x)$ for $J_2(\alpha_i) = 0$.

Here we have that

$$c_i = \frac{2}{[J_3(\alpha_i)]^2} \int_0^1 x^3 J_2(\alpha_i x) dx.$$

We can compute the integral using a substitution of $t = \alpha_i x$ and the identity in Equation (4).

$$\int_{0}^{1} x^{3} J_{2}(\alpha_{i}x) dx = \frac{1}{\alpha_{i}^{4}} \int_{0}^{\alpha_{i}} t^{3} J_{2}(t) dt
= \frac{1}{\alpha_{i}^{4}} \int_{0}^{\alpha_{i}} \frac{d}{dt} \left[t^{3} J_{3}(x) \right] dt
= \frac{1}{\alpha_{i}^{4}} \left[t^{3} J_{3}(x) \right]_{0}^{\alpha_{i}}
= \frac{1}{\alpha_{i}} J_{3}(\alpha_{i}).$$
(6)

This gives

$$c_i = \frac{2}{J_3(\alpha_i)}$$

and

$$x^{2} = \sum_{i=1}^{\infty} \frac{2}{J_{3}(\alpha_{i})} J_{2}(\alpha_{i}x), \tag{7}$$

where $J_2(\alpha_i) = 0$. In Figure 1 is a plot of this function with 100 terms.

Example 2. Consider the function f(x) = 1, $x \in [0, 2]$. Find the Fourier-Bessel series expansion in $J_0(\alpha_i x)$ for $J'_0(2\alpha_i) = 0$.

Here we have that for i > 1,

$$c_i = \frac{1}{2 \left[J_0(2\alpha_i) \right]^2} \int_0^2 x J_0(\alpha_i x) \, dx.$$

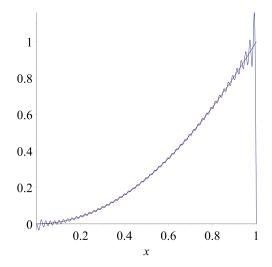


Figure 1: Plot of Fourier-Bessel series expansion for $f(x) = x^2$ using 100 terms.

We can compute the integral using a substitution of $t = \alpha_i x$ and the identity in Equation (??).

$$\int_{0}^{2} x J_{0}(\alpha_{i}x) dx = \frac{1}{\alpha_{i}^{2}} \int_{0}^{2\alpha_{i}} t J_{0}(t) dt
= \frac{1}{\alpha_{i}^{2}} \int_{0}^{2\alpha_{i}} \frac{d}{dt} [t J_{1}(x)] dt
= \frac{1}{\alpha_{i}^{2}} [t J_{1}(x)]_{0}^{2\alpha_{i}}
= \frac{2}{\alpha_{i}} J_{1}(2\alpha_{i}).$$
(8)

This gives

$$c_i = \frac{1}{2 \left[J_0(2\alpha_i) \right]^2} \frac{2}{\alpha_i} J_1(2\alpha_i) = \frac{J_1(2\alpha_i)}{\alpha_i \left[J_0(2\alpha_i) \right]^2}.$$

We note that for i = 1, we also have

$$c_1 = \frac{2}{b^2} \int_0^b x f(x) dx = \frac{1}{2} \int_0^2 x dx = 1.$$

Then, we find that

$$1 = 1 + \sum_{i=1}^{\infty} \frac{J_1(2\alpha_i)}{\alpha_i \left[J_0(2\alpha_i)\right]^2} J_0(\alpha_i x), \tag{9}$$

In Figure 2 is a plot of this function with 10 terms. where $J_0'(2\alpha_i) = 0$.

The theory and history of series involving Bessel function is given in Chapter XVIII of Watson's A Treatise on the Theory of Bessel Functions. Series of the form

$$f(x) = \sum_{m=1}^{\infty} a_m J_0(j_m x), \quad 0 < x < 1,$$

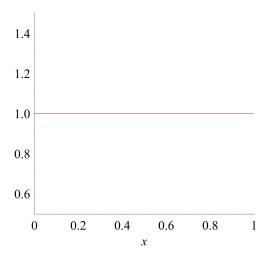


Figure 2: Plot of Fourier-Bessel series expansion for f(x) = 1 using 10 terms.

where the j_m are positive zeros of $J_0(x)$ were treated in studies by Fourier of the conduction of heat, Daniel Bernoulli's vibrating chain, and Euler's vibrating circular membrane.

This can be extended to the Fourier-Bessel series

$$f(x) = \sum_{m=1}^{\infty} a_m J_n(j_m x), \quad 0 < x < 1,$$

where the j_m are positive zeros of $J_n(x)$. Multiplying the series by $xJ_n(j_kx)$ and integrating, one has

$$\int_{0}^{1} x f(x) J_{n}(j_{k}x) dx = \sum_{m=1}^{\infty} a_{m} \int_{0}^{1} x J_{n}(j_{m}x) J_{n}(j_{k}x) dx$$
$$= a_{k} ||J_{n}(j_{k}x)||^{2} \delta_{k,m}.$$
(10)

According to Watson, the expansion coefficients can be determined from

$$2\lambda_k^2 \int_0^1 x f(x) J_n(\lambda_k x) \, dx = [(\lambda_k^2 - n^2) J_n^2(\lambda_k) + \lambda_k^2 {J_n'}^2(\lambda_k)] a_k$$

So, if $\lambda_k = j_k$ are the zeros of $J_n(x)$, then

$$||J_n(j_k x)||^2 = \frac{1}{2} J_n'^2(\lambda_k)$$

and

$$a_k = \frac{2}{J_n'^2(\lambda_k)} \int_0^1 x f(x) J_n(j_k x) dx.$$

In the more general case, one can have mixed boundary conditions, f'(1) + Af(1) = 0. Then, the λ_n 's are roots of $z^{-n}[zJ_n(z) + AJ_n(z)]$. The resulting series is called a Dini series. In the next example we derive the formula giving for $||J_n(\alpha x)||^2$.

Example 3. Show that

$$\int_0^b x [J_n(\alpha x)]^2 dx = \frac{1}{2\alpha^2} \left[\left(\alpha^2 b^2 - n^2 \right) J_n^2(\alpha b) + \alpha^2 b^2 J_n'^2(\alpha b) \right].$$

We begin with the Bessel equation (3) for different values of α . Consider the equations

$$\frac{d}{dx}\left(x\frac{dJ_n(\alpha x)}{dx}\right) + \left(\alpha^2 x - \frac{n^2}{x}\right)J_n(\alpha x) = 0.$$
(11)

$$\frac{d}{dx}\left(x\frac{dJ_n(\beta x)}{dx}\right) + \left(\beta^2 x - \frac{n^2}{x}\right)J_n(\beta x) = 0.$$
(12)

Multiply Equation (11) by $J_n(\beta x)$, Equation (12) by $J_n(\alpha x)$, subtract and rearrange to find

$$J_{n}(\beta x) \left[\frac{d}{dx} \left(x \frac{dJ_{n}(\alpha x)}{dx} \right) \right] - J_{n}(\alpha x) \left[\frac{d}{dx} \left(x \frac{dJ_{n}(\beta x)}{dx} \right) \right] = (\beta^{2} - \alpha^{2}) x J_{n}(\alpha x) J_{n}(\beta x)$$

$$\frac{d}{dx} \left[x J_{n}(\beta x) \frac{dJ_{n}(\alpha x)}{dx} - x J_{n}(\alpha x) \frac{dJ_{n}(\beta x)}{dx} \right] = (\beta^{2} - \alpha^{2}) x J_{n}(\alpha x) J_{n}(\beta x). \tag{13}$$

Integrate from x = 0 to x = b, to obtain

$$(\beta^{2} - \alpha^{2}) \int_{0}^{b} x J_{n}(\alpha x) J_{n}(\beta x) dx = \left[x J_{n}(\beta x) \frac{dJ_{n}(\alpha x)}{dx} - x J_{n}(\alpha x) \frac{dJ_{n}(\beta x)}{dx} \right]_{0}^{b}$$
$$= b J_{n}(\beta b) \frac{dJ_{n}(\alpha x)}{dx} - b J_{n}(\alpha b) \frac{dJ_{n}(\beta x)}{dx}. \tag{14}$$

So, if $\alpha \neq \beta$ and $AJ_n(\alpha b) + \alpha BJ'_n(\alpha b) = 0$, then

$$\int_0^b x J_n(\alpha x) J_n(\beta x) \, dx = 0.$$

However, what if $\alpha = \beta$? Then, we need to compute

$$L = \lim_{\alpha \to \beta} \frac{1}{\beta^2 - \alpha^2} \left[x J_n(\beta x) \frac{dJ_n(\alpha x)}{dx} - x J_n(\alpha x) \frac{dJ_n(\beta x)}{dx} \right]_0^b$$

using l'Hopital's Rule. First, we simplify the numerator using the identity

$$\frac{dJ_n(z)}{dz} = \frac{n}{z}J_n(z) - J_{n+1}(z).$$

Let $z = \alpha x$. Then,

$$\frac{dJ_n(z)}{dx} = \frac{dJ_n(z)}{dz} \frac{dz}{dx}
= \alpha \left[\frac{n}{z} J_n(z) - J_{n+1}(z) \right]
\frac{dJ_n(\alpha x)}{dx} = \frac{n}{x} J_n(\alpha x) - \alpha J_{n+1}(\alpha x).$$
(15)

So, we have

$$\left[xJ_{n}(\beta x)\frac{dJ_{n}(\alpha x)}{dx} - xJ_{n}(\alpha x)\frac{dJ_{n}(\beta x)}{dx}\right]_{0}^{b} = bJ_{n}(\beta b)\left[\frac{n}{b}J_{n}(\alpha b) - \alpha J_{n+1}(\alpha b)\right] - bJ_{n}(\alpha b)\left[\frac{n}{b}J_{n}(\beta b) - \beta J_{n+1}(\beta b)\right] = \beta bJ_{n}(\alpha b)J_{n+1}(\beta b) - \alpha bJ_{n+1}(\alpha b)J_{n}(\beta b) \tag{16}$$

We also have the identity

$$\frac{dJ_n(z)}{dz} = J_{n-1}(z) - \frac{n}{z}J_n(z).$$

Therefore, we can use the result

$$\frac{dJ_{n+1}(\alpha x)}{d\alpha} = J_n(\alpha x) - \frac{n+1}{\alpha} J_{n+1}(\alpha x)$$

to simplify the numerator in the limit:

$$\frac{d}{d\alpha} \left[\beta b J_n(\alpha b) J_{n+1}(\beta b) - \alpha b J_{n+1}(\alpha b) J_n(\beta b) \right] = \beta b J_{n+1}(\beta b) \left(\frac{n}{\alpha} J_n(\alpha b) - x J_{n+1}(\alpha b) \right) \\
- b J_{n+1}(\alpha b) J_n(\beta b) \\
- \alpha b J_n(\beta b) \left(b J_n(\alpha b) - \frac{n+1}{\alpha} J_{n+1}(\alpha b) \right), \quad (17)$$

Then the limit simplifies as

$$L = \lim_{\alpha \to \beta} \frac{\frac{d}{d\alpha} \left[\beta b J_n(\alpha b) J_{n+1}(\beta b) - \alpha b J_{n+1}(\alpha b) J_n(\beta b)\right]}{-2\alpha}$$

$$= b \frac{J_{n+1}(\beta b) \left[n J_n(\beta b) - \beta b J_{n+1}(\beta b)\right] - J_{n+1}(\beta b) J_n(\beta b) - J_n(\beta b) \left[\beta b J_n(\beta b) - (n+1) J_{n+1}(\beta b)\right]}{-2\beta}$$

$$= \frac{b}{2\beta} \left[\beta b J_n^2(\beta b) - 2n J_n(\beta b) J_{n+1}(\beta b) + \beta b J_{n+1}^2(\beta b)\right]. \tag{18}$$

Since

$$J_{n+1}(z) = \frac{n}{z}J_n(z) - J'_n(z),$$

we can write the result in terms of $J_n(\beta b)$ and $J'_n(\beta b)$ as

$$L = \frac{b}{2\beta} \left[\beta b J_n^2(\beta b) - 2n J_n(\beta b) \left(\frac{n}{\beta b} J_n(\beta b) - J_n'(\beta b) \right) + \beta b \left(\frac{n}{\beta b} J_n(\beta b) - J_n'(\beta b) \right)^2 \right]$$

$$= \frac{b}{2\beta} \left[\left(\beta b - \frac{2n^2}{\beta b} + \frac{n^2}{\beta b} \right) J_n^2(\beta b) + (2n - 2n) J_n(\beta b) J_n'(\beta b) + \beta b J_n'^2(\beta b) \right]$$

$$= \frac{1}{2\beta^2} \left[\left(\beta^2 b^2 - n^2 \right) J_n^2(\beta b) + \beta^2 b^2 J_n'^2(\beta b) \right]. \tag{19}$$

We have shown that

$$\int_{0}^{b} x [J_{n}(\alpha x)]^{2} dx = \frac{1}{2\alpha^{2}} \left[\left(\alpha^{2} b^{2} - n^{2} \right) J_{n}^{2}(\alpha b) + \alpha^{2} b^{2} J_{n}^{\prime 2}(\alpha b) \right]. \tag{20}$$